

Thermal-mechanical constraining of large scale ice flow models in Antarctica.

Eric Larour
Eric Rignot
Hélène Seroussi
Mathieu Morlighem



1. Introduction
2. Higher order inverse control methods.
3. Large scale modeling using Anisotropic Mesh Adaptation.
4. Ice flow model of Antarctica using ISSM.
5. Perspectives



1 Introduction

- Large scale modeling of Antarctica:
 - 1km resolution on Antarctica -> 20 Million elements in 2d
 - 400 million in 3d (20 vertical layers)
 - Full stokes: 1.6 billion dofs. (4 per node)
 - Cost is prohibitive.
- Constraints on bedrock friction and ice rheology:
 - Paleo runs for large scale models are hard to converge to present time.
 - Paleo runs usually do not account for full stress equilibrium (SIA).
 - A mix of paleo run and inverse control methods at present time could be necessary (similar to GCM spin up).



2 Higher order inverse control methods.

- Cost function:
$$J = \iint_{\text{Surface}} \frac{1}{2} \left\{ (u - u_{obs})^2 + (v - v_{obs})^2 \right\} dx dy$$
- We augment J with the ice flow model desired, multiplied by adjoint vectors. The model equations depend on the order modeling desired:

- Macayéal:
$$J = \iint_s \frac{1}{2} \left\{ (u - u_{obs})^2 + (v - v_{obs})^2 \right\} dx dy +$$

$$\iint_s \lambda_x(x, y) \left\{ \frac{\partial}{\partial x} \left(2\nu H \left(2 \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) \right) + \frac{\partial}{\partial y} \left(\nu H \left(\frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} \right) \right) - \rho g H \frac{\partial z_s}{\partial x} - \beta^2 u \right\} dx dy +$$

$$\iint_s \lambda_y(x, y) \left\{ \frac{\partial}{\partial y} \left(2\nu H \left(2 \frac{\partial v}{\partial y} + \frac{\partial u}{\partial x} \right) \right) + \frac{\partial}{\partial x} \left(\nu H \left(\frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} \right) \right) - \rho g H \frac{\partial z_s}{\partial y} - \beta^2 v \right\} dx dy$$



- Pattyn:
$$J = \iint_{\text{Surface}} \frac{1}{2} \left\{ (u - u_{obs})^2 + (v - v_{obs})^2 \right\} dx dy +$$

$$\iint_{\text{Volume}} \lambda_x(x, y) \left\{ \frac{\partial}{\partial x} \left(2\nu \left(2\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) \right) + \frac{\partial}{\partial y} \left(\nu \left(\frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} \right) \right) + \frac{\partial}{\partial z} \left(\nu \frac{\partial u}{\partial z} \right) - \rho g \frac{\partial z_s}{\partial x} \right\} dx dy +$$

$$\iint_{\text{Volume}} \lambda_y(x, y) \left\{ \frac{\partial}{\partial x} \left(\nu \left(\frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} \right) \right) + \frac{\partial}{\partial y} \left(2\nu \left(2\frac{\partial v}{\partial y} + \frac{\partial u}{\partial x} \right) \right) + \frac{\partial}{\partial z} \left(\nu \frac{\partial v}{\partial z} \right) - \rho g \frac{\partial z_s}{\partial y} \right\} dx dy$$

- Second part of misfit is integrated on the volume instead of the surface. Macayael is thickness integrated, Pattyn is 3d. Drag is a boundary condition for Pattyn, instead of a surface term for MacAyeal.



- Stokes:**

$$J = \iint_{\text{Surface}} \frac{1}{2} \left\{ (u - u_{obs})^2 + (v - v_{obs})^2 \right\} dx dy +$$

$$\iint_{\text{Volume}} \lambda_x(x, y) \left\{ \frac{\partial}{\partial x} \left(2\nu \frac{\partial u}{\partial x} \right) + \frac{\partial}{\partial y} \left(\nu \left(\frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} \right) \right) + \frac{\partial}{\partial z} \left(\nu \left(\frac{\partial u}{\partial z} + \frac{\partial w}{\partial x} \right) \right) - \frac{\partial P}{\partial x} \right\} dx dy +$$

$$\iint_{\text{Volume}} \lambda_y(x, y) \left\{ \frac{\partial}{\partial x} \left(\nu \left(\frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} \right) \right) + \frac{\partial}{\partial y} \left(2\nu \frac{\partial v}{\partial y} \right) + \frac{\partial}{\partial z} \left(\nu \left(\frac{\partial v}{\partial z} + \frac{\partial w}{\partial y} \right) \right) - \frac{\partial P}{\partial y} \right\} dx dy$$

$$\iint_{\text{Volume}} \lambda_z(x, y) \left\{ \frac{\partial}{\partial x} \left(\nu \left(\frac{\partial u}{\partial z} + \frac{\partial w}{\partial x} \right) \right) + \frac{\partial}{\partial y} \left(\nu \left(\frac{\partial v}{\partial z} + \frac{\partial w}{\partial y} \right) \right) + \frac{\partial}{\partial z} \left(2\nu \frac{\partial w}{\partial z} \right) - \frac{\partial P}{\partial z} - \rho g \right\} dx dy$$

$$\iint_{\text{Volume}} \lambda_p(x, y) \left\{ \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} \right\} dx dy$$
- Add vertical stress equilibrium + incompressibility equation.
 Observations misfit still integrated over surface layer.



- Misfit gradients with respect to drag coefficient:

$$u = k.N_{eff}^p \sigma_{drag}^q \quad \text{or} \quad \sigma_{drag} = \alpha^2 u^r N_{eff}^{-s} \text{ (Paterson, 1994)}$$

- Stokes:

$$\begin{aligned} \frac{\partial J}{\partial \alpha} = & -\lambda_x \left(2\alpha(v_x - v_z n_x n_z) \right) - \lambda_x \left(2\alpha(v_x - v_z n_x n_z) \right) \\ & - \lambda_z \left(2\alpha(-v_x n_x n_z - v_y n_y n_z) \right) \end{aligned}$$

- Pattyn and MacAyeal:

$$\frac{\partial J}{\partial \alpha} = -2\alpha(\lambda_x v_x + \lambda_y v_y)$$



- Thermal constraints:
 - Ice rheology inverted for ice shelves (no bedrock friction).
 - For ice sheet: inversion of drag and ice rheology is a severely underconstrained optimization problem -> solution exhibits multiple extrema.

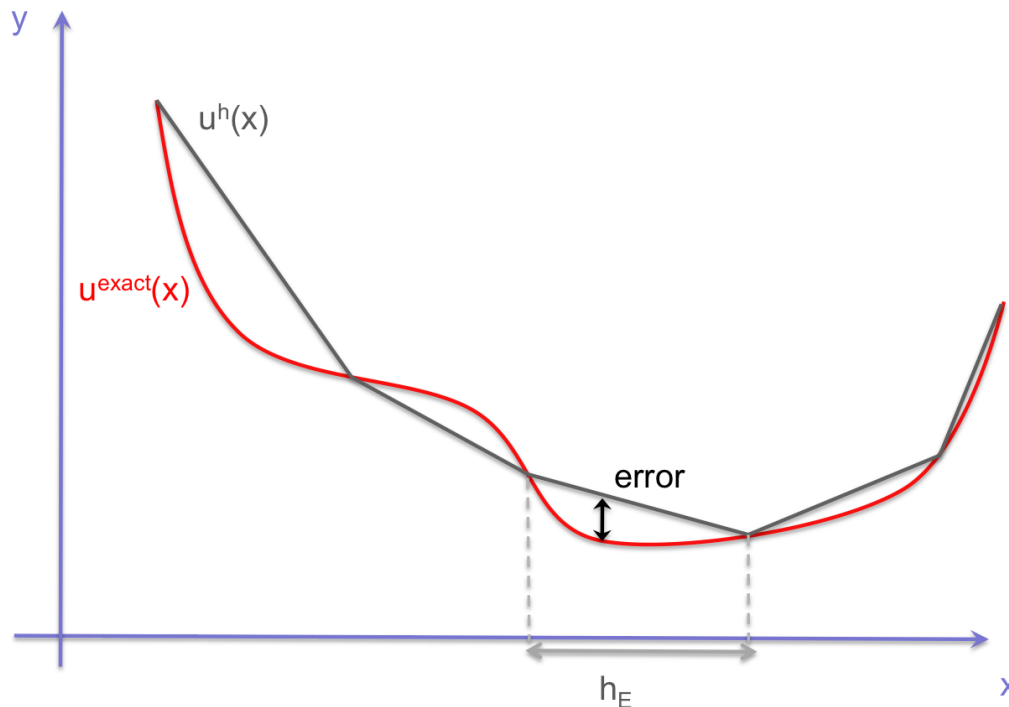
We run steady-state thermal model at each iteration of the inverse control method, so that thermally induced stresses are accounted for. Thermal model is equally computationally intensive.



3 Large scale modeling using Anisotropic Mesh Adaptation.

- If the solution $u(x)$ is approximated by $u_h(x)$, with piecewise linear interpolation, a local approximation error can be defined over an element E to be :

$$x \in [0; h_E] : error = |u^{exact}(x) - u_E^h(x)|$$





Generalized error estimate:

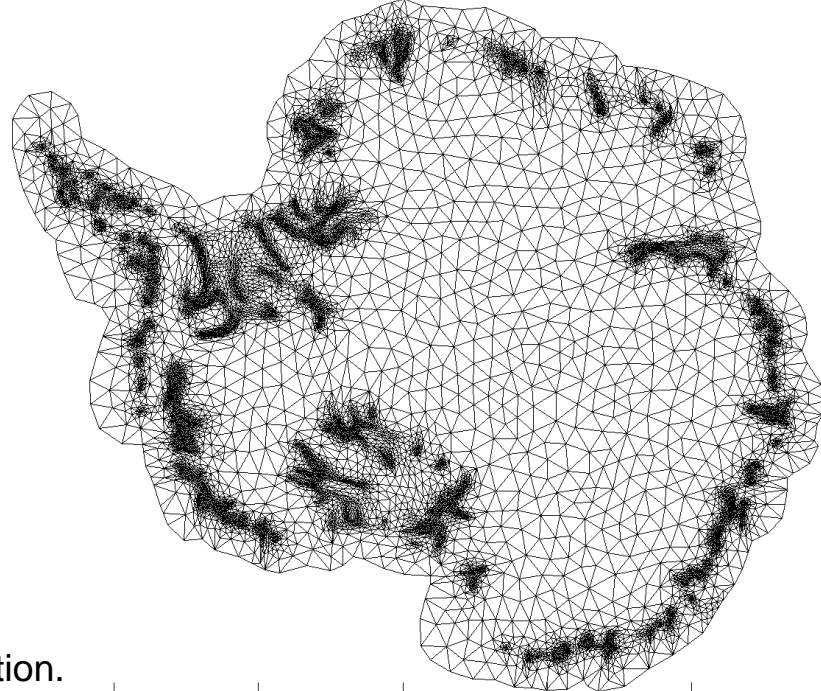
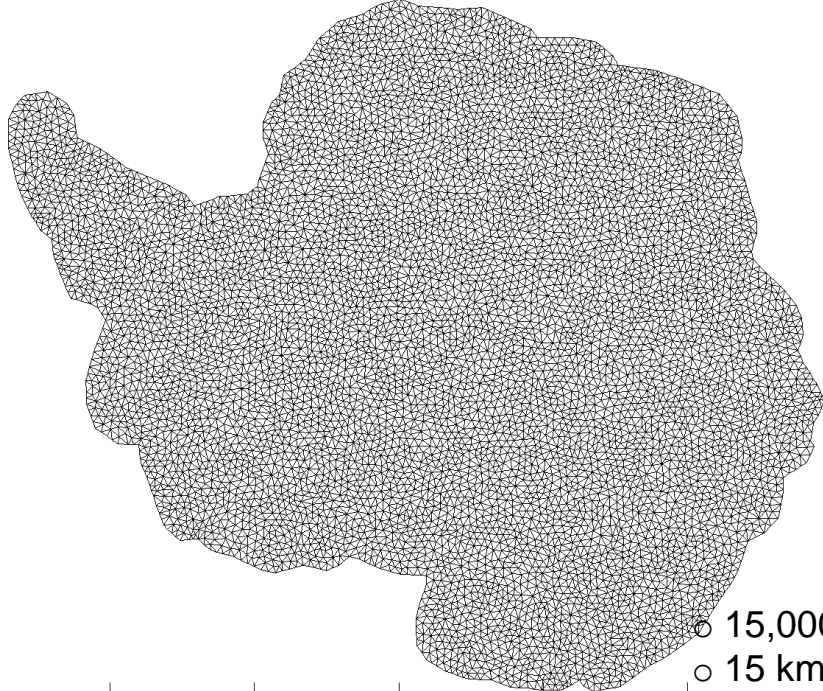
$$\left| u^{exact}(x) - u_E^h(x) \right| \leq c_d h_E^2 \sup_{(x,y) \in E} |H_u(x,y)| \quad (\text{Habashi 2000})$$

where :

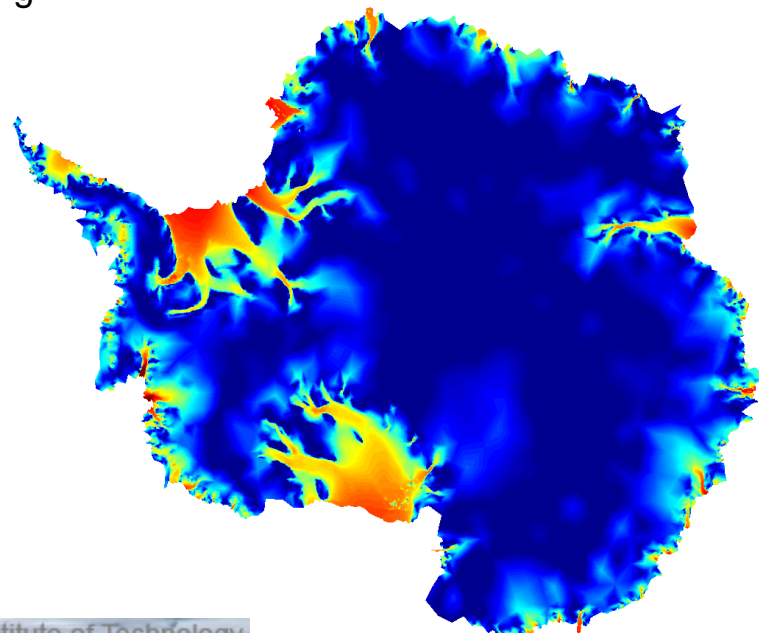
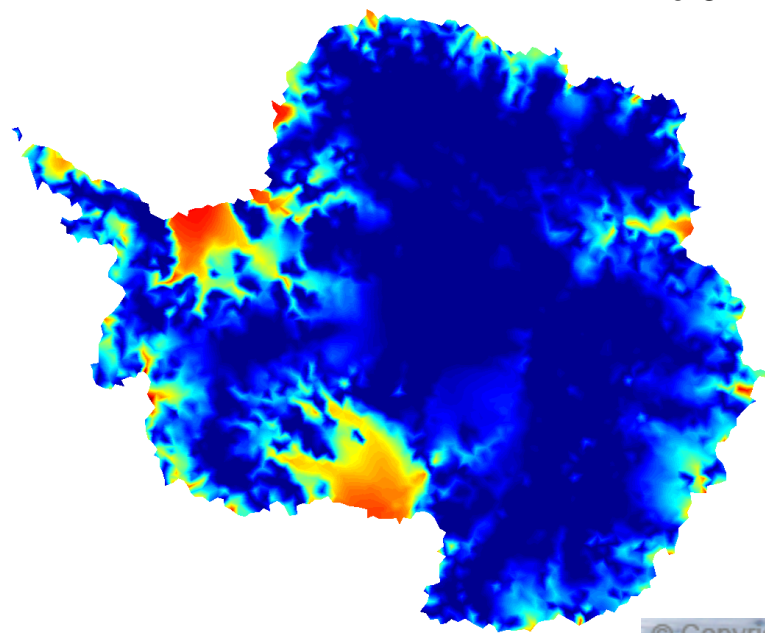
- h_E length of the element edge
- C_d constant that depends only on the space dimension (1.8 in 1d, 2.9 in 2d)
- $H_f(x; y)$ Hessian matrix of u , $|Hu(x; y)|$ its spectral norm

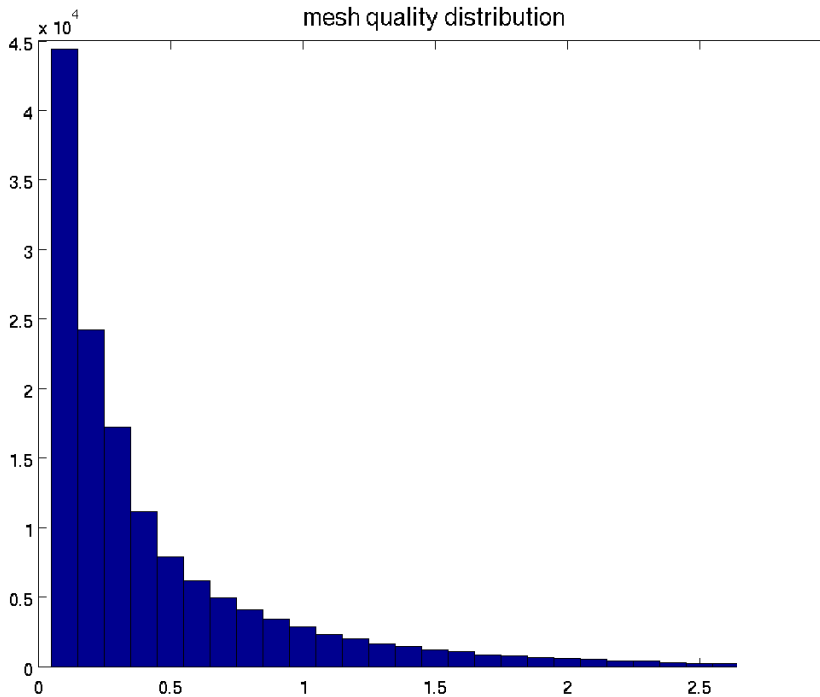
-> use Hessian matrix to minimize the error estimate, by remeshing along principal directions of Hessian matrix, according to eigenvalue magnitude.

Tool: YAMS, developed within the GAMMA research project at INRIA-Rocquencourt. Anisotropic. Pascal Frey.



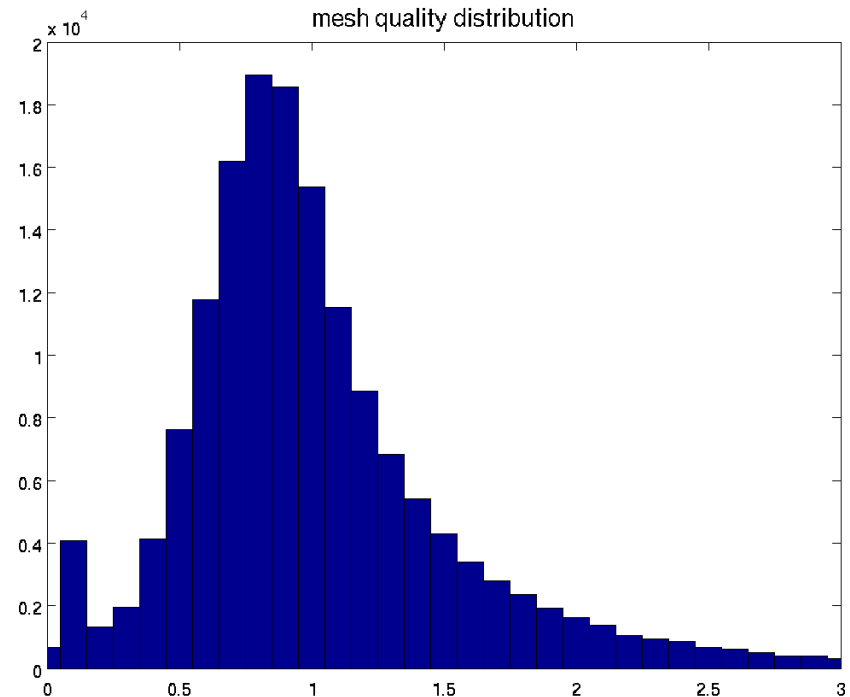
- 15,000 elements.
- 15 km initial resolution.
- 5km final resolution on icestreams.
- 3 min Yams remeshing.





In transformed error coordinates space (along Hessian directions), mesh triangles should tend to be equilateral (best capture of discretization error).

Mesh quality: measure of distortion from equilateral discretization error. Tends to 1 for equilateral triangles in error space.



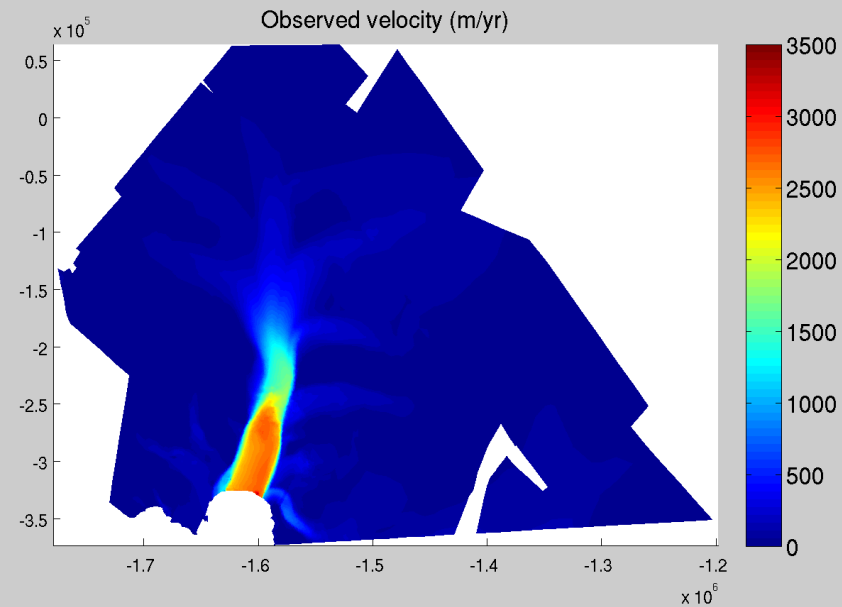
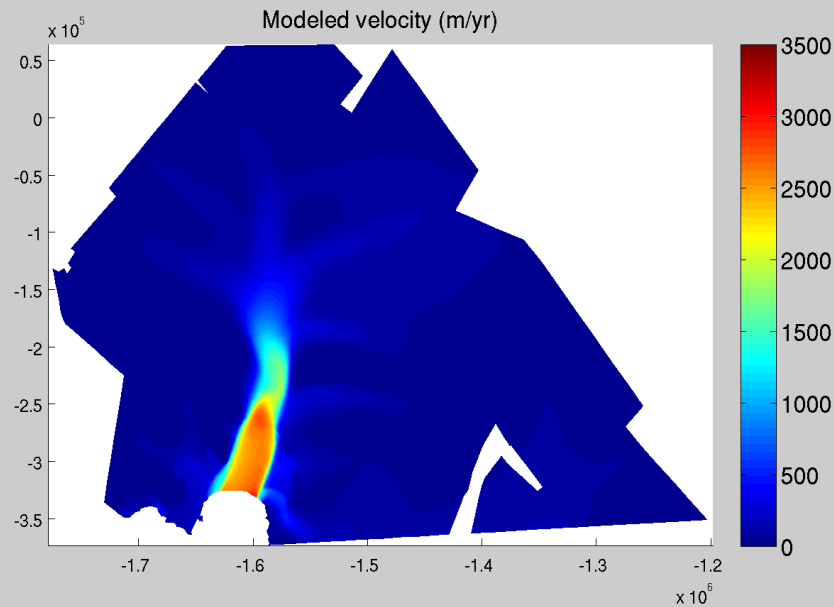


4 Ice flow model of Antarctica using ISSM.

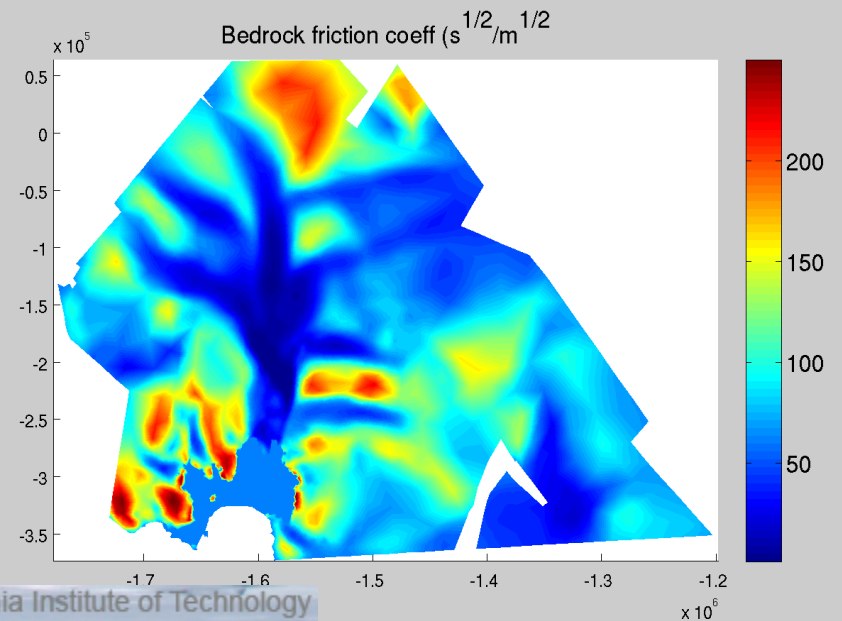
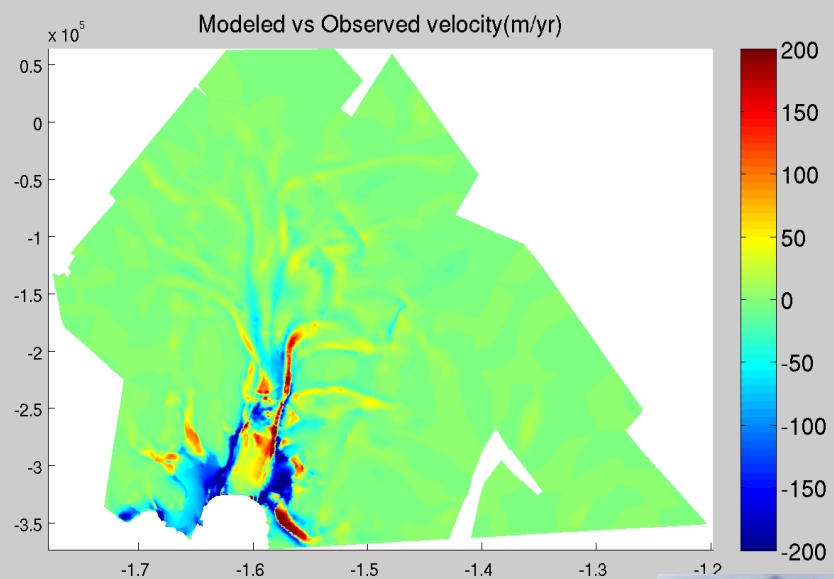
- ISSM: Ice Sheet System Model, developed by JPL's R&TD program, funded by JPL and NASA (Map09).
- Large scale model of Antarctica, using anisotropic remeshing:
 - 150,000 2d elements: MacAyeal formulation.
 - 1,200,000 3d elements (8 extrusion layers, distorted towards bedrock). Pattyn formulation.Icestreams resolved at 3km, interior of ice sheet captured at 50km.
- Diagnostic run, constrained using inverse control methods on drag:
 - Background run (40 iterations) to correctly constrain entire ice sheet.
 - Refinement on all basins (20 basins) to capture icestreams.

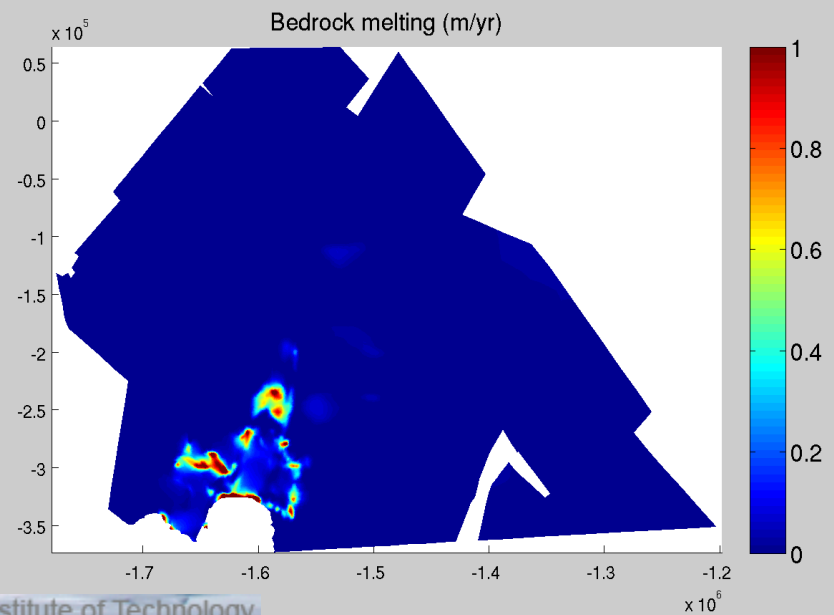
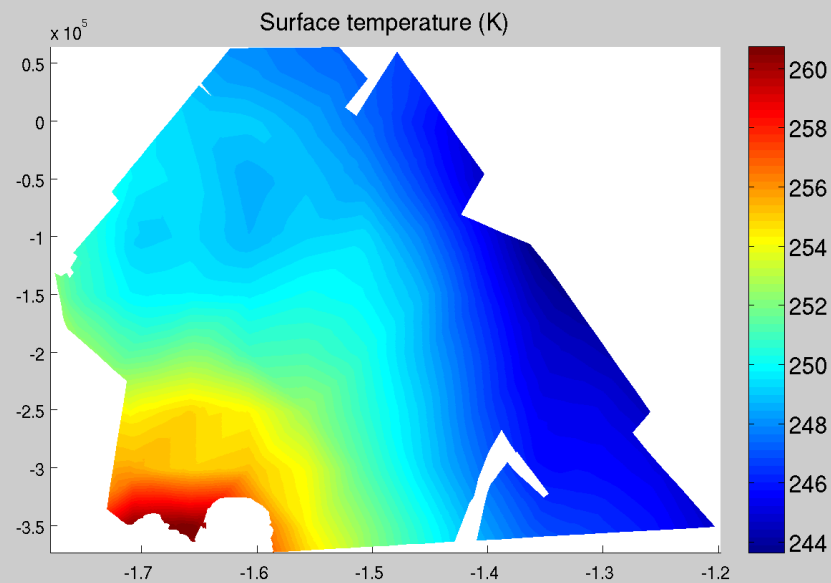
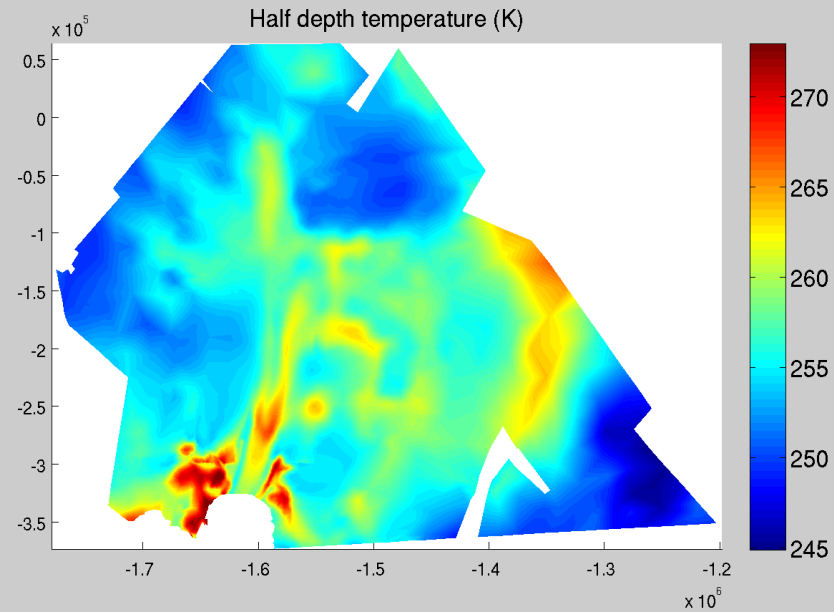
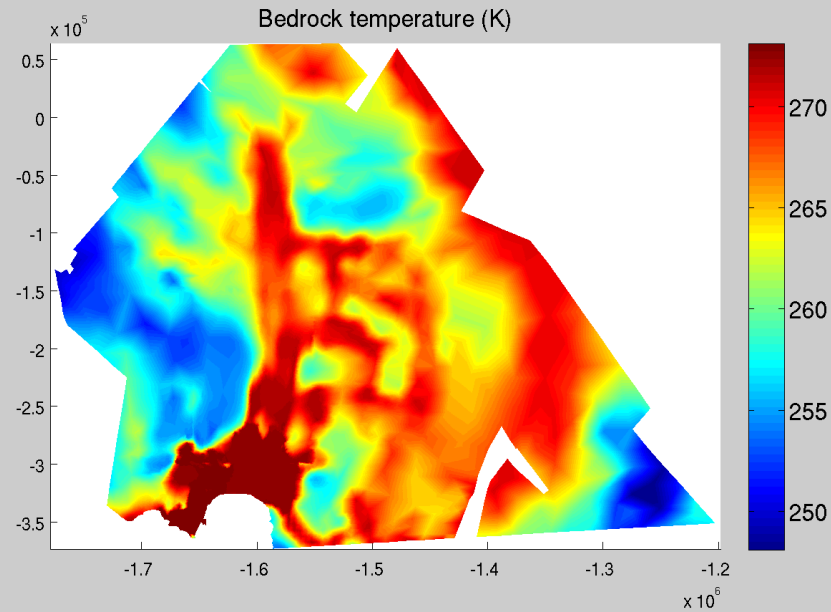


- Firn Layer: van den Broeke, M.R., Towards quantifying the contribution of the Antarctic ice sheet to global sea level change. *Journal of Physics. IV France*, 2006 (139) 170-187
- Temperatures: Giovinetto, M.B., N.M. Waters, and C.R. Bentley, Dependence of Antarctic surface mass balance on temperature, elevation and distance to open ocean, *Journal of Geophysical Research*, 1990, 95, 3517-3531
- Surface: Bamber, Jonathan L., Jose Luis Gomez-Dans, and Jennifer A. Griggs. 2009. *Antarctic 1 km Digital Elevation Model (DEM) from Combined ERS-1 Radar and ICESat Laser Satellite Altimetry*. Boulder, Colorado USA: National Snow and Ice Data Center. Digital media.
- Thickness: Lythe, M.B., D.G. Vaughan and Consortium BEDMAP, BEDMAP: A new ice thickness and subglacial topographic model of Antarctica, *Journal of Geophysical Research*, 2001, 106 (B6), 11,335-11,352
- Grounding Line, Ice Front, Ice Rises: Rignot unpublished.
- Surface velocity map: Rignot, unpublished.

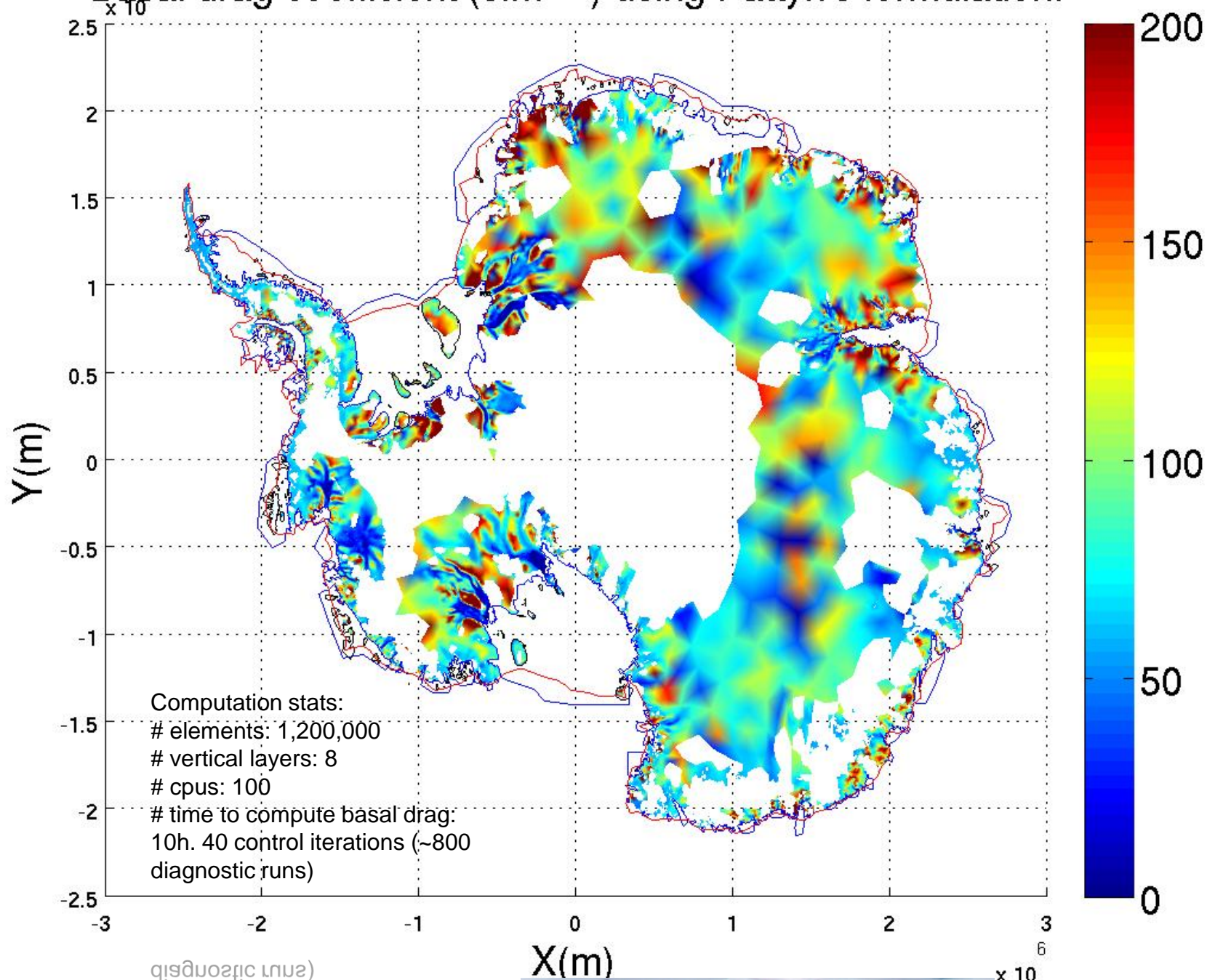


Cf: Mathieu Morlighem poster, Tuesday session.

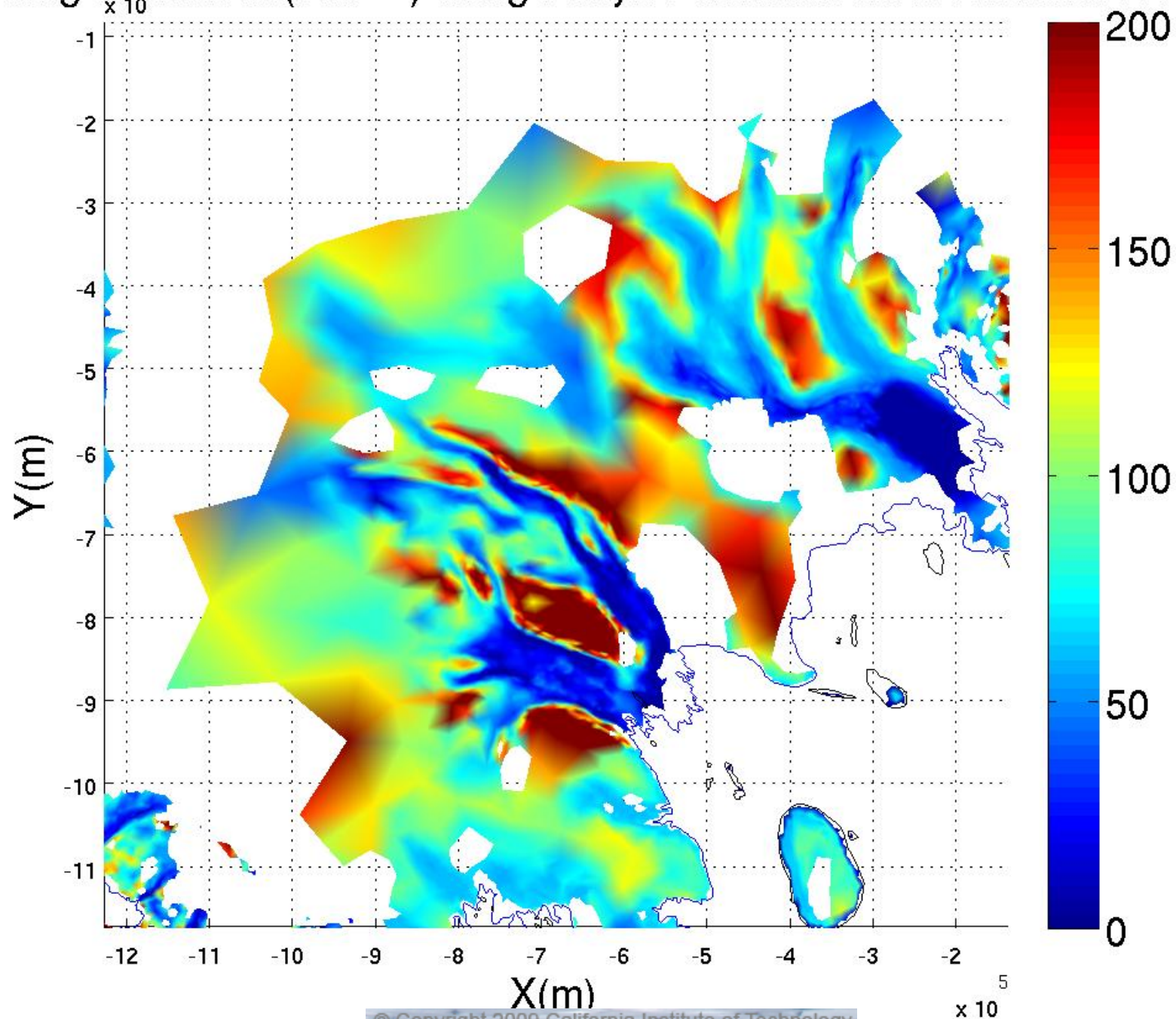




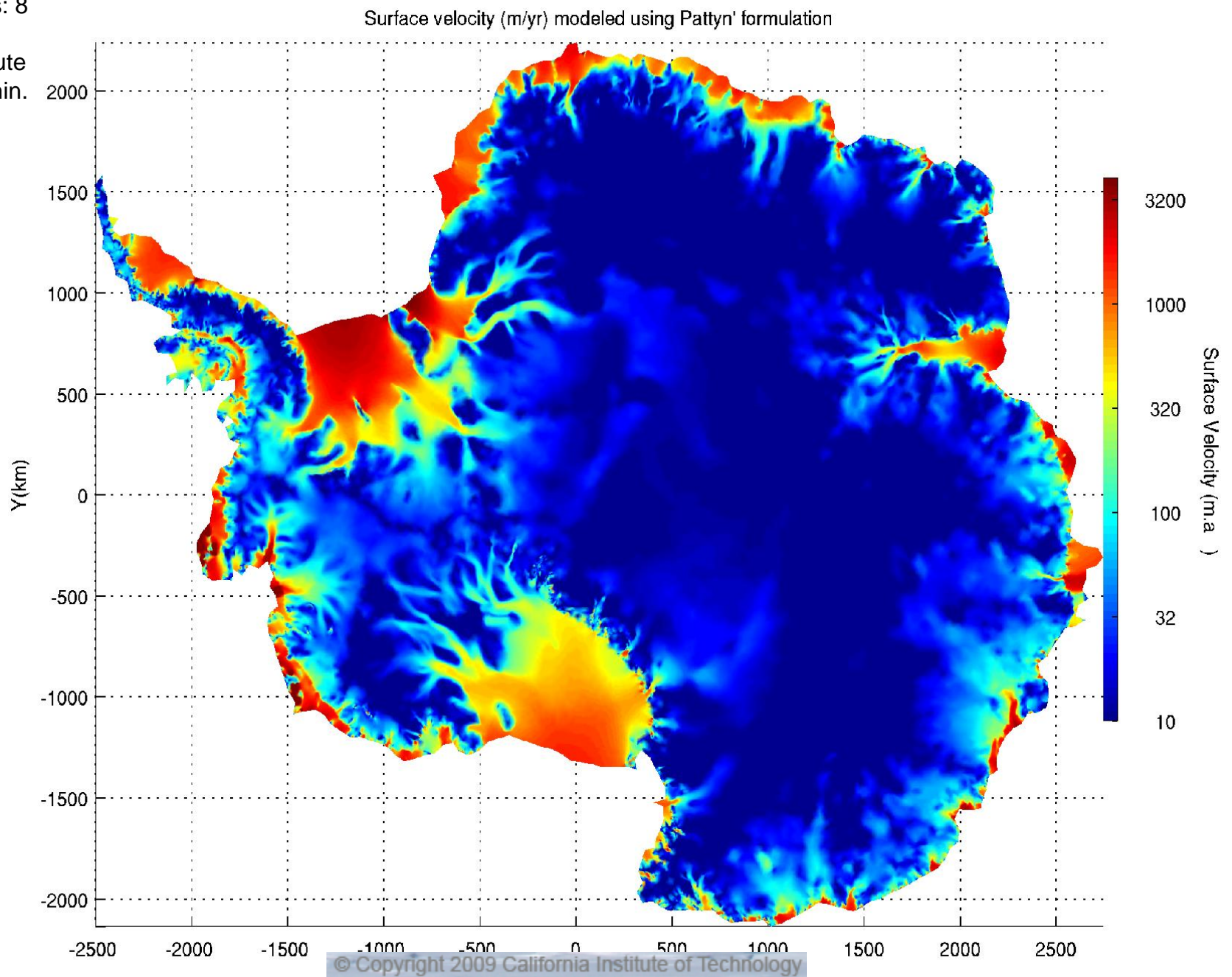
Basal drag coefficient ($\text{s.m}^{-1/2}$) using Pattyn's formulation.



Basal drag coefficient ($\text{s.m}^{-1/2}$) using Pattyn's formulation. Icestreams A->E.



Computation stats:
#elements: 1,200,000
#vertical layers: 8
#cpus: 100
#time to compute
diagnostic: 5 min.





National Aeronautics and
Space Administration

Jet Propulsion Laboratory
California Institute of Technology
Pasadena, California

5 Conclusions and perspectives.

- Higher order inverse control methods are computationally affordable, using adaptative remeshing.
- InSAR data becoming available to constrain entire continent.
- Spin ups can now combine paleo-runs with inverse control methods to constrain Antarctica ice flow.
- ISSM capable of fully constraining present day diagnostics, with assumption of thermal steady-state.
- Embedded Full-Stokes inversion computationally possible in the next couple years.
- Short term transients should be possible with full resolution models.

This work was performed at the California Institute of Technology's Jet Propulsion Laboratory under a contract with the National Aeronautics and Space Administration's Cryosphere Science Program.



THANKS!